# ab initio SCF-MO-LCAO CALCULATIONS WITH A MOLECULECALIBRATED CONTRACTED GAUSSIAN BASIS. I. 

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Molecule-calibrated contractions with Gaussian type function contraction coefficients chosen from expansion coefficients of molecular orbitals of $\mathrm{CH}_{3}^{+}$and $\mathrm{CH}_{3}^{-}$are tested. A possibility is examined to transfer these contracted functions from one ion to the another and to ethylene. Total energy with all molecules under study is lower if molecule-calibrated contracted functions are employed as a basis set instead of contracted functions derived from the examination of atomic orbitals. The effect of the molecule-calibrated contraction on energy components, orbital energies and atomic populations is also discussed.

In ab initio calculations using a Gaussian basis set all integrals necessary for the solution of Hartree-Fock-Roothaan equations ${ }^{1}$ can be evaluated analytically ${ }^{2}$ rather easily. However, one must employ a large number of Gaussian type functions (GTF), compared with that of Slater type functions used in calculations yielding results of the same accuracy ${ }^{3}$. This is a disadvantage, because the number of integrals to be calculated and employed in a SCF procedure increases with a fourth power of the number of basis functions. It is possible to overcome this difficulty by means of "contracted Gaussian type functions" (CGTF), i.e. by the adoption of a linear combination of Gaussian functions instead of individual GTF's ${ }^{4}$. A similar treatment of "lobe" Gaussian type functions was independently reported by Preuss ${ }^{5}$ and Whitten ${ }^{6}$. An actual construction of a CGTF set is equivalent to forming suitable fixed GTF linear combinations, based on the analysis of atomic orbital expansion coefficients. Examples can be found e.g. in the papers of Clementi ${ }^{4,7}$ and Salez and Veillard ${ }^{8}$. CGTF bases formed in this way have been employed in molecular calculations without any remarkable effect on the quality of results provided the contraction is well chosen ${ }^{8-10}$. Salez and Veillard ${ }^{8}$ showed that the best contracted bases for molecules are not identical with those for atoms. This is due to a deformation of atomic orbitals, in particular of the outer ones forming the chemical bonds. This should be taken into account in the choice of a contraction for molecular calculations.

Recently some authors ${ }^{10-12}$ used molecule-calibrated contracted (MCC) Gaussian basis in molecular calculations. Finding the MCC Gaussian basis set requires a SCF calculation on a small molecule using an uncontracted basis. From the examination of the molecular orbital expansion coefficients obtained the linear combinations of the GTF's are set up, the fixed coefficients of which are taken from the suitable molecular orbitals. The contracted Gaussian type
functions (MCC GTF) formed in this way are then employed as the basis for calculations of larger molecules. The choice of the fixed coefficients should be unambiguous at least with one contracted orbital for an atomic shell. From this point of view symmetrical molecules with degenerate molecular orbitals are particularly suitable. Another problem in the choice of the molecule used for calibration concerns a possibility to transfer contracted functions for an atom to another molecule containing that atom. With the experience, concerning a transfer of parameters from a small system to a larger one, achieved with localized orbitals ${ }^{13-16}$, it can be anticipated that such a transfer could be possible provided the atom lies in a similar environment and has the same valence angles and the hybrid orbitals of the same type both in a calibrated and the considered molecule.

In this paper properties of the CGTF's constructed from molecular orbitals of planar $\mathrm{CH}_{3}^{+}$and $\mathrm{CH}_{3}^{-}$ions are examined. Their transfer from the positive ion to the negative ion and vice versa and their adoption in the calculation on ethylene have been investigated.

## METHOD OF CALCULATION

The program used is an extensively modified version of the Polyatom system, QCPE 47•1 (rcf. ${ }^{17}$ ) for the IBM 7040 computer adapted to make use of contracted functions. In general a contraction does not lead to a decrease in a number of integrals over GTF's to be explicitely calculated, because each integral over a contracted function is calculated as a sum of integrals over uncontracted functions. It decreases merely the number of integrals stored on a magnetic tape and the number of linear variation parameters in the SCF calculation. However, if an integral contained two or more CGTF's which were centered on the same nucleus, then a part of the computation of all integrals over the GTF's composing these CGTF's was the same.

In the present paper a series of calculations was performed in which the bases differed only by several CGTF's. The modified program permits to use the integrals already calculated and stored on a magnetic tape in further calculations on a given molecule. Thus only integrals over new CGTF's have to be calculated.

The GTF basis set employed is of the $(7 s, 3 p / 3 s)$ type, from which the minimal contracted $[2 s, 1 p / 1 s]$ basis was formed. The latter is the same basis as that employed by Clementi in calculations on a series of molecules ${ }^{18.19}$. This basis appears to be suitable for calculations on larger molecules as it yields results of sufficient accuracy within reasonable computer times. Exponents and contraction coefficients of the GTF's for carbon and hydrogen are listed in Tables I and II. The contracted functions are not normalized, their normalization is performed before the SCF procedure starts.

Molecular-calibrated contractions were made in a way presented in the following description. First, SCF calculations using the uncontracted basis were carried out on planar $\mathrm{CH}_{3}^{+}$and $\mathrm{CH}_{3}^{-}$ions with $D_{3 \mathrm{~b}}$ symmetry setting the principal axis along the $z$ direction and putting the bond length equal to 2.025 Bohr (nuclear repulsion energy 9.744227 a.u.). The calculations proceeded until the difference in electronic energies in two succesive iterations was less than $10^{-7}$ a.u. A linear combination of five s-type GTF's possessing the highest exponents with coefficients taken from the $1 a_{1}^{\prime} \mathrm{MO}$
represents the $1 s$ MCC GTF, a linear combination of two s-type GTF's with the lowest exponents with coefficients taken from the $2 a_{1}^{\prime} \mathrm{MO}$ forms the 2 s MCC GTF, a linear combination of three $p$-type GTF's with coefficients taken from the $1 e^{\prime}$ and

Table I
Atomic Contraction Coefficients and MCC Coefficients for Carbon Atom Derived from Molecular Orbitals of $\mathrm{CH}_{3}^{+}$and $\mathrm{CH}_{3}^{--}$Ions

| Exponent | Orbital representation | Coefficient |  |  |
| :---: | :---: | :---: | :---: | :---: |
|  |  | atomic contraction | calibrated molecule |  |
|  |  |  | $\mathrm{CH}_{3}^{+}$ | $\mathrm{CH}_{3}^{-}$ |
| 391.445 | $1 s$ | 0.02220 | 0.02219 | 0.02218 |
| 64.7358 |  | 0.13285 | 0.13279 | 0.13277 |
| 16.2247 |  | 0.38435 | 0.38419 | 0.38404 |
| $5 \cdot 33460$ |  | 0.45798 | 0.45702 | 0.45653 |
| 2.00995 |  | 0.15441 | 0.15539 | 0.15582 |
| 0.502323 | $2 s$ | 0.56673 | 0.55775 | 0.41758 |
| 0.155139 |  | 0.55692 | 0.32580 | 0.30068 |
| 4.31613 | $p_{\text {x,y }}$ | $0 \cdot 10845$ | 0.08182 | 0.06359 |
| 0.873682 |  | 0.46116 | 0.37017 | 0.27279 |
| 0.202860 |  | 0.63043 | 0.36763 | 0.25492 |
| 4.31613 | $p_{\text {z }}$ | $0 \cdot 10845$ | 0.08862 | 0.09256 |
| 0.873682 |  | 0.46116 | 0.35370 | 0.34048 |
| 0.202860 |  | 0.63043 | 0.73542 | 0.74515 |

Table II
Atomic Contraction Coefficients and MCC Coefficients for Hydrogen Atom Derived from Molecular Orbitals of $\mathrm{CH}_{3}^{+}$and $\mathrm{CH}_{3}^{-}$Ions

| Exponent | Type | atomic contraction | Coefficient |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  | calibrated molecule |  |  |  |  |  |
|  |  |  | $\begin{aligned} & \mathrm{CH}_{3}^{+} \\ & \left(2 a_{1}^{\prime}\right) \end{aligned}$ | $\begin{aligned} & \mathrm{CH}_{3}^{+} \\ & \left(1 e^{\prime}\right) \end{aligned}$ | $\begin{gathered} \mathrm{CH}_{3}^{+} \\ \left(\mathrm{mix}^{a}\right) \end{gathered}$ | $\begin{aligned} & \mathrm{CH}_{3}^{-} \\ & \left(2 a_{1}^{\prime}\right) \end{aligned}$ | $\begin{aligned} & \mathrm{CH}_{3}^{-} \\ & \left(I e^{\prime}\right) \end{aligned}$ | $\underset{\left(\mathrm{mix}^{d}\right)}{\mathrm{CH}_{3}^{-}}$ |
| 0.151374 | $s$ | 0.64767 | 0.01889 | 0.06325 | $0 \cdot 19950$ | 0.08725 | 0.27141 | $1 \cdot 10060$ |
| 0.681277 | $s$ | 0.40789 | 0.14078 | 0.23870 | 1.00000 | 0.13773 | 0.23463 | 1.00000 |
| 4.50037 | $s$ | 0.07048 | 0.02808 | 0.04466 | 0.19306 | 0.02941 | 0.04623 | $0 \cdot 20494$ |

[^0]$1 a_{2}^{\prime \prime}$ MO's, respectively, represents the $p_{x, y}$ and $p_{z}$ MCC GTF's for carbon. For hydrogen the contraction was made by a linear combination of the GTF's with coefficients taken from both the $2 a_{1}^{\prime}$ and $1 e^{\prime}$. MO's of $\mathrm{CH}_{3}^{+}$and $\mathrm{CH}_{3}^{-}$ions. Moreover, a "mixture" of linear combinations with coefficients from the $2 a_{1}^{\prime}$ and $1 e^{\prime}$ MO's was made for both ions, because the GTF's on hydrogen do not enter these linear combinations with the same weight. The ratio of coefficients in this contraction is an average value of their ratios in the $2 a_{1}^{\prime}$ and $1 e^{\prime}$ MO's. The ratio $1: 2$ of coefficients was also tested, but this contraction appeared to be less successful. A similar treatment was employed by Hoyland ${ }^{11}$ in estimates of barriers of internal rotation in alkanes,

Table III
Basis Sets for Carbon and Hydrogen Atoms

| Number of the basis set | Contraction |
| :---: | :---: |
| 1 | No contraction |
| 2 | $1 s(\mathrm{~A}) 2 s(\mathrm{~A}) p_{\mathrm{x}, \mathrm{y}}(\mathrm{A}) p_{\mathrm{z}}(\mathrm{A}) / h(\mathrm{~A})$ |
| 3 | $1 s\left(\mathrm{CH}_{3}^{+}\right) 2 s(\mathrm{~A}) p_{\mathrm{x}, \mathrm{y}}(\mathrm{A}) p_{z}(\mathrm{~A}) / h(\mathrm{~A})$ |
| 4 | $1 s(\mathrm{~A}) 2 s\left(\mathrm{CH}_{3}^{+}\right) p_{x, y}(\mathrm{~A}) p_{\mathrm{z}}(\mathrm{A}) / h(\mathrm{~A})$ |
| 5 | $1 s\left(\mathrm{CH}_{3}^{+}\right) 2 s\left(\mathrm{CH}_{3}^{+}\right) p_{\mathrm{x}, \mathrm{y}}(\mathrm{A}) p_{\mathrm{z}}(\mathrm{A}) / h(\mathrm{~A})$ |
| 6 | $1 s(\mathrm{~A}) 2 s(\mathrm{~A}) p_{\mathrm{x}, \mathrm{y}}\left(\mathrm{CH}_{3}^{+}\right) p_{\mathrm{z}}(\mathrm{A}) / h(\mathrm{~A})$ |
| 7 | $1 s(\mathrm{~A}) 2 s(\mathrm{~A}) p_{\mathrm{x}, \mathrm{y}}(\mathrm{A}) p_{\mathrm{z}}\left(\mathrm{CH}_{3}^{+}\right) / h(\mathrm{~A})$ |
| 8 | $1 s\left(\mathrm{CH}_{3}^{+}\right) 2 s\left(\mathrm{CH}_{3}^{+}\right) p_{\mathrm{x}, \mathrm{y}}\left(\mathrm{CH}_{3}^{+}\right) p_{z}\left(\mathrm{CH}_{3}^{+}\right) / h(\mathrm{~A})$ |
| 9 | $1 s\left(\mathrm{CH}_{3}^{-}\right) 2 s(\mathrm{~A}) p_{x, y}(\mathrm{~A}) p_{z}(\mathrm{~A}) / h(\mathrm{~A})$ |
| 10 | $1 s(\mathrm{~A}) 2 s\left(\mathrm{CH}_{3}^{-}\right) p_{x, y}(\mathrm{~A}) p_{\mathrm{z}}(\mathrm{A}) / h(\mathrm{~A})$ |
| 11 | $1 s\left(\mathrm{CH}_{3}^{-}\right) 2 s\left(\mathrm{CH}_{3}^{-}\right) p_{\mathrm{x}, \mathrm{y}}(\mathrm{A}) p_{\mathrm{z}}(\mathrm{A}) / h(\mathrm{~A})$ |
| 12 | $1 s(\mathrm{~A}) 2 s(\mathrm{~A}) p_{x, y}\left(\mathrm{CH}_{3}^{-}\right) p_{\mathrm{z}}(\mathrm{A}) / h(\mathrm{~A})$ |
| 13 | $1 s\left(\mathrm{CH}_{3}^{-}\right) 2 s\left(\mathrm{CH}_{3}^{-}\right) p_{x, y}\left(\mathrm{CH}_{3}^{-}\right) p_{z}(\mathrm{~A}) / h(\mathrm{~A})$ |
| 14 | $1 s(\mathrm{~A}) 2 s(\mathrm{~A}) p_{\mathrm{x}, \mathrm{y}}(\mathrm{A}) p_{\mathrm{z}}\left(\mathrm{CH}_{3}^{-}\right) / h(\mathrm{~A})$ |
| 15 | $1 s\left(\mathrm{CH}_{3}^{-}\right) 2 s\left(\mathrm{CH}_{3}^{-}\right) p_{\mathrm{x}, \mathrm{y}}\left(\mathrm{CH}_{3}^{-}\right) p_{\mathrm{z}}\left(\mathrm{CH}_{3}^{-}\right) / h(\mathrm{~A})$ |
| 16 | $1 s(\mathrm{~A}) 2 s\left(\mathrm{CH}_{3}^{+}\right) p_{x, y}\left(\mathrm{CH}_{3}^{-}\right) p_{z}\left(\mathrm{CH}_{3}^{+}\right) / h(\mathrm{~A})$ |
| 17 | $1 s(\mathrm{~A}) 2 s\left(\mathrm{CH}_{3}^{+}\right) p_{\chi, y}\left(\mathrm{CH}_{3}^{-}\right) p_{\mathrm{z}}(\mathrm{A}) / h(\mathrm{~A})$ |
| 18 | $16^{a} / h\left(\mathrm{CH}_{3}^{+} 2 a_{1}^{\prime}\right)$ |
| 19 | $16^{a} / h\left(\mathrm{CH}_{3}^{+} 1 e^{\prime}\right)$ |
| 20 | $16^{a} / h\left(\mathrm{CH}_{3}^{-} 22 a_{1}^{\prime}\right)$ |
| 21 | $\left.16^{a} / h\left(\mathrm{CH}_{3}^{-}\right) 1 e^{\prime}\right)$ |
| 22 | $8^{a} / h\left(\mathrm{CH}_{3}^{+}\right.$mix $)$ |
| 23 | $15^{a} / h\left(\mathrm{CH}_{3}^{-}\right.$mix) |
| 24 | $17^{a} / h\left(\mathrm{CH}_{3}^{+}\right.$mix $)$ |
| 25 | $17^{a} / h\left(\mathrm{CH}_{3}^{-} \mathrm{mix}\right)$ |

[^1]constructing the MCC GTF basis from results of calculations on methane using an uncontracted basis set and also by Schulman, Moskowitz and Hollister ${ }^{10}$ in calculations on ethylene.

For judging the quality of the MCC GTF basis set a series of calculations was performed in which the atomic contraction was substituted stepwise by the molecular one for the $1 s, 2 s, p_{\mathrm{x}, \mathrm{y}}, p_{\mathrm{z}}$ orbitals and their combinations. 25 various basis sets formed in this way are summarized in Table III. As an example the basis 18 is expressed in its analytical form; the carbon CGTF's are as follows

$$
\begin{aligned}
\chi_{1 \mathrm{~s}} & =0.02220(391.445)+0.13285(64.7358)+0.38435(16.2247) \\
& +0.45798(5.33460)+0.15441(2.00995) \\
\chi_{2 \mathrm{~s}} & =0.55775(0.502323)+0.32580(0.155139), \\
\chi_{p_{\mathrm{x} .5}} & =0.06359(4.31613)+0.27279(0.873682)+0.25492(0.20286), \\
\chi_{\mathbf{p}_{z}} & =0.08862(4.31613)+0.35370(0.873682)+0.73542(0.20286),
\end{aligned}
$$

Table IV
Total Energy, Energy Components and Orbital Energies for the $\mathrm{CH}_{3}^{+}$Ion (in a. u.)

| Number <br> of the <br> basis set <br> from | $E_{\text {tot. }}$ | $E_{\text {pot }}$. | $E_{\text {kin }}$ | $-V / 2 T$ |
| :--- | :--- | :--- | :--- | :--- | | One- |
| :---: |
| Table III |


| 1 | -39.156744 | -78.2155 | 39.0588 | 1.0012 | -107.0721 |
| ---: | :--- | :--- | :--- | :--- | :--- |
| 2 | -38.916853 | -77.3130 | 38.3962 | 1.0068 | -105.2081 |
| 3 | -38.916814 | -77.3051 | 38.3883 | 1.0069 | -105.1992 |
| 5 | -38.938523 | -77.5431 | 38.6046 | 1.0043 | -105.5323 |
| 6 | -38.958167 | -77.5524 | 38.5942 | 1.0047 | -105.6518 |
| 8 | -38.978799 | -77.7719 | 38.7931 | 1.0024 | -105.9604 |
| 9 | -38.916782 | -77.3025 | 38.3857 | 1.0069 | -105.1962 |
| 11 | -38.931445 | -77.4568 | 38.5254 | 1.0053 | -105.4135 |
| 12 | -38.963337 | -77.6217 | 38.6584 | 1.0039 | -105.7615 |
| 15 | -38.977115 | -77.7569 | 38.7798 | 1.0025 | -105.9537 |
| 16 | -38.983788 | -77.8455 | 38.8617 | 1.0016 | -106.0736 |
| 18 | -39.114301 | -79.0724 | 39.9581 | 0.9894 | -108.2372 |
| 19 | -39.121426 | -78.8042 | 39.6828 | 0.9929 | -107.7421 |
| 20 | -39.106076 | -78.3846 | 39.2785 | 0.9978 | -106.9627 |
| 21 | -39.044777 | -77.9895 | 38.9447 | 1.0013 | -106.2842 |
| 22 | -39.121706 | -78.8217 | 39.7000 | 0.9927 | -107.7886 |
| 23 | -39.047440 | -77.9244 | 38.8770 | 1.0022 | -106.2009 |

and the hydrogen CGTF is

$$
\chi_{\mathrm{h}}=0.01889(0.151374)+0.14078(0.681277)+0.02808(4.50037) .
$$

The numbers in parentheses are the GTF exponents. The expansion coefficients of the GTF's are taken from Tables I and II.

## RESULTS AND DISCUSSION

## $\mathrm{CH}_{3}^{+}$and $\mathrm{CH}_{3}^{-}$Ions

Both ions were calculated assuming the same geometry regardless of whether the MCC GTF or uncontracted basis set was used. The results are presented in Tables IV - VII. The molecule-calibrated contraction for the carbon 1 s orbital (bases 3, 9) gives practically the same results as the atomic contraction (basis 2 ), which indicates that the deformation of the $1 s$ orbital in the molecules studied is unimportant. The molecule-calibrated contraction for the $2 s, p_{x, y}, p_{z}$ carbon orbitals appears to be

Table IV
(Continued)

superior to the atomic contraction (basis 2) considering the total energy. The respective decrease in the total energy is roughly additive, e.g. with $\mathrm{CH}_{3}^{+}$employing the basis 8, where the MCC for carbon is made from the MO's of $\mathrm{CH}_{3}^{-}$the decrease in energy amounts to 0.061949 a.u. while if the MCC is separately carried out for the $2 s$ and $p_{x, y}$ orbitals using the bases 5 and 6 , the sum of the respective energy decreases amounts to 0.062984 a.u. (The total energy of $\mathrm{CH}_{3}^{+}$is independent on the way in which the contraction of the $p_{\mathrm{z}}$ orbital is done.)
For the $2 s$ carbon orbital the MCC formed from the MO's of $\mathrm{CH}_{3}^{+}$is better for both ions, for the $p_{x, y}$ orbitals the MCC formed from the MO's of $\mathrm{CH}_{3}^{-}$was better, for the $p_{\mathrm{z}}$ orbital of $\mathrm{CH}_{3}^{-}$the MCC formed from the MO's of $\mathrm{CH}_{3}^{+}$was better. In general the best contraction of carbon orbitals with both ions is the basis 16 , where each contraction for any orbital was that which gave better results concerning the total energy. Using this basis the total energy, energy components and orbital energies of $\mathrm{CH}_{3}^{+}$are more close to values yielded by an uncontracted basis than those

Table V
Total Energy, Energy Components and Orbital Energies for the $\mathrm{CH}_{3}^{-}$Ion (in a. u.)

| Number <br> of the <br> basis set <br> from | $E_{\text {tot. }}$ | $E_{\text {pot. }}$ | $E_{\text {kin }}$. | $-V / 2 T$ |
| :---: | :---: | :---: | :---: | :---: | | One- |
| :---: |
| elcctron |
| Table III |


| 1 | -39.369184 | -78.9099 | 39.5407 | 0.9978 | -114.3053 |
| ---: | ---: | ---: | :--- | :--- | :--- |
| 2 | -39.272661 | -78.8385 | 39.5658 | 0.9963 | -113.9885 |
| 3 | -39.272655 | -78.8308 | 39.5582 | 0.9964 | -113.9803 |
| 5 | -39.279658 | -78.9638 | 39.6841 | 0.9949 | -114.1784 |
| 6 | -39.285824 | -78.9665 | 39.6807 | 0.9950 | -114.2283 |
| 7 | -39.289261 | -78.5048 | 39.2155 | 1.0009 | -113.2795 |
| 8 | -39.313679 | -78.7798 | 39.4661 | 0.9981 | -113.7443 |
| 9 | -39.272643 | -78.8282 | 39.5556 | 0.9964 | -113.9773 |
| 11 | -39.277575 | -78.9153 | 39.6377 | 0.9954 | -114.1075 |
| 12 | -39.287561 | -79.0040 | 39.7164 | 0.9946 | -114.2845 |
| 13 | -39.292298 | -79.0774 | 39.7851 | 0.9938 | -114.3981 |
| 14 | -39.288712 | -78.4958 | 39.2071 | 1.0010 | -113.2458 |
| 15 | -39.313358 | -78.7618 | 39.4485 | 0.9983 | -113.7016 |
| 16 | -39.315937 | -78.8267 | 39.5108 | 0.9975 | -113.8129 |
| 18 | -39.180182 | -80.6472 | 41.4670 | 0.9724 | -117.7138 |
| 19 | -39.237967 | -80.3249 | 41.0869 | 0.9775 | -117.0140 |
| 20 | -39.313282 | -79.7790 | 40.4657 | 0.9858 | -115.7910 |
| 21 | -39.340301 | -79.1487 | 39.8084 | 0.9941 | -114.4603 |
| 22 | -39.222161 | -80.4022 | 41.1800 | 0.9762 | -117.2208 |
| 23 | -39.340820 | -79.1429 | 39.8021 | 0.9942 | -114.4674 |

obtained by means of the basis 2 . However, the results of the population analysis are better with the basis 2 ; here the respective deviations from the results of the calculation using the uncontracted basis set fall in opposite directions depending on whether the basis 2 or 16 is used. Using the former, carbon is more negative than using the uncontracted basis, while using the latter, it is more positive in agreement with the results obtained with a larger basis $(10 s, 6 p / 4 s)$ (gross populations for carbon and hydrogen are 5.870 and 0.710 , respectively ${ }^{20}$.

The calculated total energy of $\mathrm{CH}_{3}^{-}$using the basis 16 is 0.043276 a.u. lower compared to the value obtained using the basis 2 . Most other quantities studied are however closer in value to those computed with an uncontracted basis if the basis 2 is employed.

Concerning the hydrogen MCC GTF's constructed from the MO's of $\mathrm{CH}_{3}^{+}$, it is seen in Table V, that their transfer to $\mathrm{CH}_{3}^{-}$does not give good results (bases 18,19 , 22), MCC GTF's for hydrogens of $\mathrm{CH}_{3}^{-}$are much more diffuse than those of $\mathrm{CH}_{4}^{+}$.

Table V (continued)

| Two- <br> electron <br> potential | Orbital energies |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | $1 a_{1}^{\prime}$ | $2 a_{1}^{\prime}$ | $1 e^{\prime}$ | $1 a_{2}^{\prime \prime}$ | $3 a_{1}^{\prime}$ |
|  |  |  |  |  |  |
| 25.6512 | -10.8128 | -0.5468 | -0.2166 | 0.0616 | 0.5144 |
| 25.4058 | -10.8296 | -0.5400 | -0.2296 | 0.0230 | 0.8325 |
| 25.4052 | -10.8305 | -0.5401 | -0.2298 | 0.0227 | 0.8323 |
| 25.4704 | -10.8143 | -0.5415 | -0.2259 | 0.0282 | 0.8517 |
| 25.5176 | -10.8010 | -0.5310 | -0.2282 | 0.0322 | 0.8422 |
| 25.0305 | -10.9311 | -0.5700 | -0.2533 | 0.0064 | 0.8053 |
| 25.2203 | -10.8821 | -0.5615 | -0.2478 | 0.0202 | 0.8345 |
| 25.4048 | -10.8307 | -0.5401 | -0.2298 | 0.0227 | 0.8323 |
| 25.4480 | -10.8187 | -0.5409 | -0.2270 | 0.0267 | 0.8442 |
| 25.5363 | -10.7946 | -0.5298 | -0.2282 | 0.0331 | 0.8434 |
| 25.5764 | -10.7847 | -0.5308 | -0.2257 | 0.0366 | 0.8553 |
| 25.0057 | -10.9361 | -0.5724 | -0.2555 | 0.0058 | 0.8033 |
| 25.1955 | -10.8854 | -0.5619 | -0.2513 | 0.0187 | 0.8266 |
| 25.2420 | -10.8745 | -0.5600 | -0.2478 | 0.0210 | 0.8361 |
| 27.3224 | -10.4585 | -0.3857 | -0.0758 | 0.1947 | 1.3436 |
| 26.9449 | -10.5412 | -0.4209 | -0.1107 | 0.1649 | 1.2629 |
| 26.2678 | -10.6808 | -0.4813 | -0.1710 | 0.1092 | 1.1160 |
| 25.5674 | -10.8176 | -0.5370 | -0.2269 | 0.0498 | 0.9309 |
| 27.0744 | -10.5151 | -0.4084 | -0.0995 | 0.1764 | 1.2987 |
| 25.5802 | -10.8190 | -0.5358 | -0.2267 | 0.0517 | 0.9445 |
|  |  |  |  |  |  |

With $\mathrm{CH}_{3}^{+}$the best results were achieved with the basis 22 , which is entirely constructed from the MO's of $\mathrm{CH}_{3}^{+}$. All quantities studied are close in value to the results obtained with the uncontracted basis; in this respect they are considerably better than the values computed with the basis 2 .

The lowest total energy of $\mathrm{CH}_{3}^{-}$was achieved with the basis 23 , which is an analogue to the basis 22 for $\mathrm{CH}_{3}^{+}$, but in this case the results are not so unambiguous: some of them are better, some are worse than those computed with the basis 2, atthough the decrease in the total energy of 0.068159 a.u. is considerable.

## Ethylene

The calculations were performed assuming the molecular geometry chosen by Moskowitz and coworkers ${ }^{10.21}$, i.e. the $\mathrm{C}-\mathrm{C}$ and $\mathrm{C}-\mathrm{H}$ bond lengths 2.55 and 2.02 Bohr, respectively, and the HCH angle $120^{\circ}$. The nuclear repulsion energy amounts to 33.402244 a.u. The results are summarized in Table VIII (total energy, energy components and orbital energies) and in Table IX (population analysis). The computational costs dictated the use of the calculation of Moskowitz and Harrison ${ }^{21}$ with a ( $7 s, 3 p / 2 s$ ) basis set as a reference calculation with an uncontracted basis set.

Table VI
Population Analysis for the $\mathrm{CH}_{3}^{+}$Ion

| Number of the basis set from Table III | Net |  | Gross |  | Overlap |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | C | H | C | H |  | - . |
| 1 | 5.044 | 0.316 | 6.069 | 0.643 | 0.684 |  |
| 2 | 5.136 | 0.266 | 6.224 | 0.592 | 0.725 |  |
| 3 | $5 \cdot 136$ | 0.266 | 6.224 | 0.592 | 0.725 |  |
| 5 | 4.949 | 0.295 | 6.074 | 0.642 | 0.750 |  |
| 6 | 4.856 | 0.354 | 5.980 | 0.673 | 0.749 |  |
| 8 | 4.678 | 0.385 | 5.833 | 0.722 | 0.770 |  |
| 9 | $5 \cdot 137$ | 0.266 | 6.224 | $0 \cdot 592$ | 0.725 |  |
| 11 | 5.020 | 0.283 | 6.133 | 0.622 | $0.742$ |  |
| 12 | 4.789 | 0.378 | 5.918 | 0.694 | $0.753$ |  |
| 15 | 4.680 | 0.397 | 5.830 | 0.723 | 0.766 |  |
| 16 | 4.613 | 0.409 | $5.772$ | $0.743$ | 0.772 |  |
| 18 | 5.004 | 0.322 | 6.026 | 0.658 | 0.681 |  |
| 19 | 4.844 | 0.347 | 5.914 | 0.695 | 0.713 |  |
| 20 | 4.661 | 0.323 | 5.788 | 0.737 | 0.752 |  |
| 21 | 4.575 | 0.410 | 5.734 | 0.755 | 0.772 |  |
| 22 | 4.970 | 0.320 | 6.014 | 0.662 | 0.696 |  |
| 23 | 4.636 | 0.396 | 5.786 | 0.738 | 0.766 |  |

As the basis of that calculation differs from the basis used in the present paper in GTF exponents and by a lower number of the hydrogen GTF's, we considered also another reference calculation in which a larger basis ( $9 s, 5 p / 3 s$ ) was employed ${ }^{10}$.

From the Table VIII it is apparent that the $p_{z}$ MCC GTF's calibrated on $\mathrm{CH}_{3}^{+}$ and $\mathrm{CH}_{3}^{-}$(bases 7 and 14 , respectively) cannot be used for ethylene, as expected. $\dot{A}$ possibly way to construct a MCC $p_{\pi}$ orbital is that of exploiting the GTF coefficients of the $1 b_{14}$ MO obtained from a calculation on ethylene using an uncontracted basis set. However even the carbon atomic $p_{z}$ orbital alone appears to be a good function for description of a $\pi$ bond ${ }^{10}$.

The MCC functions for $2 s$ and $p_{x, y}$ orbitals constructed from the MO's of $\mathrm{CH}_{3}^{+}$ and $\mathrm{CH}_{3}^{-}$are very good when employed in the calculation on ethylene. With the 2 s orbital the contraction based on the MO's of $\mathrm{CH}_{3}^{+}$is slightly superior, while with $p_{\mathrm{x}, \mathrm{y}}$ orbitals that based on the MO 's of $\mathrm{CH}_{3}^{-}$is better, similarly as in the case of $\mathrm{CH}_{3}^{+}$ and $\mathrm{CH}_{3}^{-}$ions. Combining these, an optimum contraction is obtained (basis 17), which gives a total energy 0.089599 a.u. lower than the calculation using an atomic

Table VII
Population Analysis for the $\mathrm{CH}_{3}^{-}$Ion

Table VIII
Total Energy, Energy Components and Orbital Energies for Ethylene (in a.u.)

| Number of the <br> basis set from <br> Table III | $E_{\text {tot. }}$ | $E_{\text {pot. }}$ | $E_{\text {kin }}$ | $-V / 2 T$ | One-electron <br> potential | Two-electron <br> potential |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $1^{a}$ | -77.800181 | - | - | - | - | - |
| 2 | -77.630942 | -154.4664 | 76.8355 | 1.0052 | -244.6744 | 56.8057 |
| 4 | -77.659217 | -154.8354 | 77.1762 | 1.0031 | -245.2335 | 56.9958 |
| 6 | -77.688232 | -154.8454 | 77.1572 | 1.0034 | -245.4398 | 57.1921 |
| 7 | -77.602572 | -154.1551 | 76.5525 | 1.0069 | -243.9881 | 56.4308 |
| 8 | -77.692341 | -154.9030 | 77.2107 | 1.0031 | -245.3201 | 57.0148 |
| 10 | -77.650159 | -154.7032 | 77.0530 | 1.0039 | -245.0344 | 56.9290 |
| 12 | -77.695793 | -154.9647 | 77.2689 | 1.0028 | -245.6355 | 57.2686 |
| 14 | -77.598605 | -154.1441 | 76.5454 | 1.0069 | -243.9509 | 56.4047 |
| 15 | -77.688463 | -154.8769 | 77.1884 | 1.0032 | -245.2798 | 57.0007 |
| 17 | -77.720541 | -155.3244 | 77.6039 | 1.0007 | -246.1835 | 57.4568 |
| 24 | -77.762637 | -155.2222 | 77.4596 | 1.0020 | -248.2861 | 59.6616 |
| 25 | -77.790950 | -155.7399 | 77.9490 | 0.9990 | -246.9903 | 57.8481 |
| $b$ | -78.0062 | - | - | - | - | - |

Table VIII

| Number of the basis set from Table III | Orbital energies |  |  |  |  |  | Orbital energies |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $1 a_{\mathrm{g}}$ | $1 b_{3 u}$ | $2 a_{8}$ | $2 b_{3 u}$ | $1 b_{2 \mathrm{u}}$ | $3 a_{\mathrm{g}}$ | $1 b_{1 g}$ | $1 b_{1 \mathrm{u}}$ | $1 b_{2 \mathrm{~g}}$ |
| $1^{a}$ | $-11.2539$ | $-11 \cdot 2526$ | $-1.0584$ | $-0.8067$ | $-0.6604$ | $-0.5829$ | $-0.5174$ | $-0.3814$ | $0 \cdot 1518$ |
| 2 | $-11.4088$ | $-11.4074$ | $-1.0947$ | $-0.8353$ | $-0.7026$ | $-0.6379$ | $-0.5636$ | $-0.4635$ | $0 \cdot 1226$ |
| 4 | $-11.3821$ | $-11.3807$ | $-1.0992$ | $-0.8353$ | $-0.6948$ | $-0.6291$ | $-0.5570$ | $-0.4546$ | 0.1319 |
| 6 | $-11.3558$ | $-11.3543$ | $-1.0755$ | $-0.8236$ | $-0.6997$ | $-0.6374$ | $-0.5571$ | $-0.4459$ | $0 \cdot 1430$ |
| 7 | $-11.4592$ | $-11.4579$ | $-1 \cdot 1146$ | $-0.8511$ | $-0.7181$ | $-0.6569$ | $-0.5778$ | $-0.4515$ | 0.0897 |
| 8 | $-11.3780$ | $-11.3766$ | $-1.0989$ | $-0.8387$ | $-0.7073$ | $-0.6480$ | $-0.5644$ | $-0.4280$ | 0.1155 |
| 10 | $-11 \cdot 3915$ | $-11 \cdot 3901$ | $-1.0980$ | $-0.8354$ | $-0.6975$ | $-0.6322$ | $-0.5593$ | $-0.4577$ | $0 \cdot 1287$ |
| 12 | $-11 \cdot 3429$ | -11.3414 | $-1.0726$ | $-0.8218$ | $-0.6991$ | $-0.6373$ | $-0.5562$ | $-0.4435$ | 0.1461 |
| 14 | $-11.4620$ | $-11.4607$ | $-1 \cdot 1162$ | $-0.8524$ | $-0.7195$ | $-0.6585$ | $-0.5790$ | $-0.4498$ | 0.0883 |
| 15 | $-11 \cdot 3770$ | $-11.3755$ | $-1.0962$ | $-0.8381$ | $-0.7110$ | $-0.6527$ | $-0.5670$ | $-0.4274$ | $0 \cdot 1736$ |
| 17 | $-11.3165$ | $-11 \cdot 3150$ | $-1.0761$ | $-0.8212$ | $-0.6912$ | $-0.6285$ | $-0.5498$ | $-0.4348$ | $0 \cdot 1551$ |
| 24 | $-11.0921$ | $-11.0906$ | $-0.9596$ | $-0.7274$ | $-0.6014$ | $-0.5079$ | $-0.4525$ | $-0.3202$ | 0.2858 |
| 25 | $-11 \cdot 2822$ | $-11.2809$ | $-1.0569$ | $-0.8083$ | $-0.6800$ | $-0.6109$ | -0.5393 | $-0.4141$ | 0.1791 |
| $b$ | $-11.2395$ | $-11.2379$ | -1.0397 | $-0.7959$ | $-0.6549$ | $-0.5812$ | $-0.5145$ | $-0.3736$ | $0 \cdot 1436$ |

[^2]contraction. Furthermore, the virial ratio 1.0007 is considerably better and also all orbital energies are closer to the values obtained with the basis 1 or the ( $9 s, 5 p / 3 s$ ) basis, which are near to the experimental values of ionization potentials ${ }^{22}$.

Also with ethylene a rough additivity in decrease of the total energy is observed, e.g. the sum of decreases in bases 4 and 12 amounts to 0.093126 a.u., compared to the total decrease of $0.089599 \mathrm{a} . \mathrm{u}$. using the basis 17 .

Considering the MCC also for hydrogen the total energy is further lowered: with the basis 25 its value is -77.790950 a.u. which is rather close to the value -77.800181 a.u. yielded by the uncontracted basis set. However it should be kept in mind that the basis 1 contains only 2 GTF's for hydrogen. When compared with the results of the calculation using the basis set 17 , orbital energies calculated with the basis 25 are even closer to those obtained with uncontracted bases with exception of the $1 b_{2_{8}}$ MO, the energy of which is too high.
In light of the successful results achieved with the MCC GTF basis set for the total and orbital energies, less satisfactory population analysis results appear to be somewhat surprising. Perhaps this can partly be assigned to the "unequilibrated" uncontracted basis sets, because a number of GTF's on hydrogen is known to affect considerably the charge distribution. From this it follows that the population analysis results are strongly dependent on the contraction employed and therefore their interpretation must be taken with some caution.

## Table IX

Population Analysis for Ethylene

| Number of the basis set from. Table III | Net |  | Gross |  | Overlap |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | C | H | C | H | C-C | $\mathrm{C}-\mathrm{H}$ |
| $1^{a}$ |  |  | $6 \cdot 330$ | 0.835 | 0.943 | 0.839 |
| 2 | 5.026 | $0 \cdot 464$ | 6.403 | 0.798 | 1.247 | 0.847 |
| 4 | 4.914 | 0.499 | 6.309 | 0.945 | $1 \cdot 251$ | 0.847 |
| 6 | 4.855 | 0.570 | 6.240 | 0.880 | 1.306 | 0.829 |
| 7 | 5.071 | 0.434 | 6.465 | 0.767 | 1.292 | 0.839 |
| 8 | 4.779 | 0.573 | 6.207 | 0.897 | 1.358 | 0.826 |
| 10 | 4.956 | 0.485 | 6.364 | 0.827 | 1.253 | 0.848 |
| 12 | 4.819 | 0.598 | $6 \cdot 200$ | 0.900 | 1.313 | 0.822 |
| 14 | 5.075 | 0.432 | 6.470 | 0.765 | 1.295 | 0.839 |
| 15 | 4.786 | 0.854 | 6.208 | 0.896 | 1.372 | 0.821 |
| 17 | 4.711 | 0.633 | 6.108 | 0.946 | 1.314 | 0.820 |
| 24 | 5.096 | 0.436 | 6.419 | 0.790 | 1.209 | 0.751 |
| 25 | 4.739 | 0.596 | 6-132 | 0.934 | -1.287 | 0.817 |

The ( $7 s, 2 p / 2 s$ ) basis set, ref. ${ }^{21}$.

## CONCLUSIONS

The results obtained can be characterized as follows:

1. A considerable decrease in the total energy was observed if the MCC GTF basis set was used instead of the CGTF's with coefficients chosen according to calculations for atoms. This energy difference amounts to 0.204853 a.u. with $\mathrm{CH}_{3}^{+}, 0.068159$ a.u. with $\mathrm{CH}_{3}^{-}$, and 0.160008 a.u. with ethylene. 2. A transfer of a MCC GTF basis set from one molecule to the another appears to be possible. With functions for hydrogen this transfer may be encountered with some difficulties. 3. Besides the total energy the quality of the MCC GTF basis set employed must also be judged by means of other quantities such as orbital energies, energy components and population analysis. 4. A rough additivity of corrections to the total energy was found, when the MCC GTF's for carbon $2 s, p_{x, y}$ and $p_{z}$ orbitals were stepwise introduced into the basis set. That means that a change in contraction coefficients in one orbital does not affect considerably a quality of contraction in another orbital. Whitman and Hornback ${ }^{23}$ found with the first period atoms that the optimum s-type GTF's exponents are not sensitive to the choice of the p-type GTF basis set and vice versa. From this and from the possibility to transfer a MCC GTF basis set from one molecule to the another it appears that a choice of molecule-calibrated GTF exponents should be possible. Some efforts have already been made in this direction ${ }^{12}$. This could be made in a following way provided the exponents for the $1 s$ function need not be optimized. First, in calculations on $\mathrm{CH}_{3}^{+}$and $\mathrm{CH}_{3}^{-}$all orbitals with exception of the $2 s$ would be contracted using the contraction coefficients chosen on the basis of calculations on atoms. The $2 s$ GTF exponents would be optimized by multiplying with a scaling parameter ${ }^{24}$. These optimized $2 s$ GTF's would first be contracted with the aid of contraction coefficients chosen from the optimum SCF calculation on the molecule and then after renormalization used in optimization of the $p_{x, y}$ GTF exponents. The exponents of the uncontracted $p_{x, y}$ GTF's would be optimized in a similar way and then the optimized GTF's would be contracted. The same treatment would be carrjed out to optimize the $p_{z}$ GTF exponents. This procedure would be repeated until convergency would be reached. The work in this direction is in progress.
[^3]
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Translated by P. Cársky.

Note added in proof: In the recent paper by Schulman, Hornback and Moskowitz ${ }^{25}$ the calculations on the ethylene with the $[5 s, 3 p, 2 d / 2 s, 1 p]$ basis set have been reported. With this basis set the gross population for carbon is $6 \cdot 215$. The gross population for carbon with our best MCC GTF basis for ethylene (basis 25 ) agrees better with this value then the gross popilation obtained with the CGTF basis optimum for isolated atoms (basis 2).


[^0]:    ${ }^{a}$ Mixture of CGTF's with coefficients from $2 a_{1}^{\prime}$ and $1 e^{\prime}$ MO's in the ratio $1: 1$.

[^1]:    ${ }^{a}$ Number of the basis set referring to the contraction on carbon atom.

[^2]:    ${ }^{6}$ The $(7 s, 3 p / 2 s)$ basis set, ref. ${ }^{21}$; ${ }^{b}$ The $(9 s, 5 p / 3 s)$ basis set, ref. ${ }^{10}$.

[^3]:    A modification of the Polyatom program for the IBM 7040 computer was made with help of Dr P. Kaiser of the Computing Laboratory, SSR Ministry of Planned Economy. His assistance is gratefully acknowledged.

